

DERIVATION OF VORTICITY TENDENCY EQUATION

vorticity: $\vec{\omega} = \nabla \times \vec{v}$

$$\vec{\omega} = (\partial_y w - \partial_z v) \hat{i} - (\partial_x w - \partial_z u) \hat{j} + (\partial_x v - \partial_y u) \hat{k}$$

Vertical component: $\xi \equiv (\partial_x v - \partial_y u) \hat{k}$

Want a time tendency (prognostic) eq. for ξ .

- ξ is a spatial derivative of \vec{v}

- Momentum equations are time tendency of \vec{v}

- time tendency of $\xi \rightarrow$ we need to take a spatial derivative of the \vec{v} tendency equations (horizontal momentum)!

Step 1: rewrite horizontal advection terms:

$$u: u \partial_x u + v \partial_y u = -v \xi + \frac{1}{2} [\partial_x (u^2 + v^2)]$$

$$= -v (\partial_x v - \partial_y u) + \frac{1}{2} [\partial_x (u^2 + v^2)]$$

u, v depend on x, y ; so

$$u \partial_x u + v \partial_y u = -v (\partial_x v - \partial_y u) + \frac{1}{2} (2(u \partial_x u + v \partial_x v))$$

$$= -v (\cancel{\partial_x v} - \partial_y u) + u \partial_x u + v \cancel{\partial_x v}$$

$$u \partial_x u + v \partial_y u = +v \partial_y u + u \partial_x u$$

similarly,

$$v: u \partial_x v + v \partial_y v = u \xi + \frac{1}{2} [\partial_y (u^2 + v^2)]$$

Now, we have horizontal advection in terms of ξ !

2. Substitute results from 1. into u & v mom^m eqs.

$$\begin{aligned} \Rightarrow u: \partial_t u - v\xi - v\eta + w\partial_z u &= -\frac{1}{\rho} \partial_x P - \frac{1}{2} \partial_x (u^2 + v^2) + \overline{X} \\ \Rightarrow v: \partial_t v + u\xi + u\eta + w\partial_z v &= -\frac{1}{\rho} \partial_y P - \frac{1}{2} \partial_y (u^2 + v^2) + \overline{Y} \end{aligned}$$

} Conservative forces
(i.e. friction
viscous forces)

3. Get ξ tendency equation

- ξ is horizontal derivative of \vec{v}
- momentum equations are \vec{v} tendency
- so, for ξ tendency, we want horiz. derivatives of momentum equations:

$$\partial_x \textcircled{2} - \partial_y \textcircled{1}$$

$$\partial_x \textcircled{2}: \partial_{x,t} v + \partial_x [u(\xi + \eta)] + \partial_x (w\partial_z v) = \partial_x \left[-\frac{1}{\rho} \partial_y P \right] - \frac{1}{2} \partial_{xy} (u^2 + v^2) + \partial_x \overline{Y}$$

$$\text{LHS: } \partial_{x,t} v + (\xi + \eta) \partial_x u + u \partial_x (\xi + \eta) + w \partial_{x,z} v + \partial_x w \partial_z v$$

$$\text{RHS: } -\frac{1}{\rho} \partial_{x,y} P + \partial_y P \partial_x \left(-\frac{1}{\rho} \right) - \frac{1}{2} \left[\partial_{xy} (u^2 + v^2) \right] + \partial_x \overline{Y}$$

$$\partial_y \textcircled{1}: \partial_{y,t} u - \partial_y [v(\xi + \eta)] + \partial_y (w\partial_z u) = \partial_y \left[-\frac{1}{\rho} \partial_x P \right] - \frac{1}{2} \partial_{yx} (u^2 + v^2) + \partial_y \overline{X}$$

$$\text{LHS: } \partial_{y,t} u - (\xi + \eta) \partial_y v - v \partial_y (\xi + \eta) + w \partial_{y,z} u + \partial_y w \partial_z u$$

$$\text{RHS: } -\frac{1}{\rho} \partial_{y,x} P + \partial_x P \partial_y \left(-\frac{1}{\rho} \right) - \frac{1}{2} \left[\partial_{yx} (u^2 + v^2) \right] + \partial_y \overline{X}$$

$\partial_x \textcircled{2} - \partial_y \textcircled{1}$, combine LHS & RHS

$$\begin{aligned} \text{LHS: } \partial_{x,t} v - \partial_{y,t} u + (\xi + \eta) (\partial_x u + \partial_y v) + u \partial_x (\xi + \eta) + v \partial_y (\xi + \eta) + w (\partial_{x,z} v - \partial_{y,z} u) \\ + \partial_x w \partial_z v - \partial_y w \partial_z u \end{aligned}$$

Recall, we can rewrite derivatives:

$$\partial_{x,t} v \Leftrightarrow \partial_{t,x} v \quad \left(\frac{\partial}{\partial x} \left(\frac{\partial v}{\partial t} \right) \Leftrightarrow \frac{\partial}{\partial t} \left(\frac{\partial v}{\partial x} \right) \right)$$

So, we can rewrite the LHS further:

$$\text{HS: } \partial_t(\partial_x v - \partial_y u) + u \partial_x(\xi + f) + v \partial_y(\xi + f) + w[\partial_z(\partial_x v - \partial_y u)] + \dots$$

$$+ (\xi + f)(\partial_x u + \partial_y v) + \partial_x w \partial_z v - \partial_y w \partial_z u$$

Now, we can substitute ξ in, and rewrite:

$$\partial_t \xi + u \partial_x(\xi + f) + v \partial_y(\xi + f) + w \partial_z \xi + (\xi + f)(\partial_x u + \partial_y v) + \dots$$

$$+ \partial_x w \partial_z v - \partial_y w \partial_z u$$

$$\frac{D}{Dt} \equiv \partial_t + u \partial_x + v \partial_y + w \partial_z, \text{ so:}$$

$$\frac{D}{Dt}(\xi + f) + (\xi + f)(\partial_x u + \partial_y v) + \partial_x w \partial_z v - \partial_y w \partial_z u: \text{ LHS}$$

↑ since f does not depend on t or z , we can put it inside $\frac{D}{Dt}$

Now, we look at the RHS:

$$\text{RHS: } \frac{1}{\rho}(\cancel{\partial_{x,y} p} - \cancel{\partial_{y,x} p}) + \partial_y p \partial_x(\frac{1}{\rho}) - \partial_x p \partial_y(\frac{1}{\rho}) - \frac{1}{2}(\cancel{\partial_{x,y}(u^2 + v^2)} - \cancel{\partial_{y,x}(u^2 + v^2)})$$

$$+ \partial_x Y - \partial_y X$$

4. Combine LHS: RHS

$$\frac{D}{Dt}(\xi + f) + (\xi + f)(\partial_x u + \partial_y v) + \partial_x w \partial_z v - \partial_y w \partial_z u = \partial_y p \partial_x(\frac{1}{\rho}) - \partial_x p \partial_y(\frac{1}{\rho}) + \partial_x Y - \partial_y X$$

Rearrange one more time:

$$\frac{D}{Dt}(\xi + f) = \underbrace{-(\xi + f)(\partial_x u + \partial_y v)}_{(1)} + \underbrace{[\partial_x u \partial_y w - \partial_x v \partial_y w]}_{(2)} + \underbrace{[\frac{1}{\rho^2}(\partial_y p \partial_x \rho - \partial_x p \partial_y \rho)]}_{(3)} + \underbrace{(\partial_x Y - \partial_y X)}_{(4)}$$

- 1) stretching / horizontal divergence
- 2) tilting (by vertical shear of u, v or horizontal shear of w)
- 3) solenoidal / baroclinic term (vertical component of $\frac{(\nabla p \times \nabla \rho)}{\rho^2}$)
- 4) nonconservative terms (e.g. friction, viscosity)

Conservation of absolute vorticity

$$\frac{D}{Dt}(\zeta + f) = 0 \quad \text{when:}$$

- incompressible fluid (ρ is 0)
- Conservative (χ is 0)
- barotropic (σ is 0)
- 2-dimensional (2 is 0)